## Bohdan Karczewski*

## Inverse Diffraction of the Electromagnetic Waves

Quite a large amount of work has recently been done on the inverse diffraction of wave fields [ 1,2 , 3, 4, 5, 6, 7]. However, the attention of research workers was almost entirely focused on the scalar wave fields. It is the aim of the present paper to give the basic formulae for the inverse diffraction of the vector electromagnetic waves. The electromagnetic theory of diffraction formulated by R. K. Luneburg [8] is taken as a starting point of our considerations. According to [8] the monochromatic, diffracted, electromagnetic field in the half-space $z>0$ may by described by the following formulae:

$$
\begin{gathered}
E_{x}\left(x_{1}, y_{1}, z_{1}\right)=\frac{1}{2 \pi} \int E_{x}\left(x_{0}, y_{0}, z_{0}\right) \frac{\partial}{\partial z_{0}}\left(\frac{e^{i k r}}{r}\right) d x_{0} d y_{0}, \\
E_{y}\left(x_{1}, y_{1}, z_{1}\right)=\frac{1}{2 \pi} \int E_{y}\left(x_{0}, y_{0}, z_{0}\right) \frac{\partial}{\partial z_{0}}\left(\frac{e^{i k r}}{r}\right) d x_{0} d y_{0},(1) \\
E_{z}\left(x_{1}, y_{1}, z_{1}\right)=\frac{1}{2 \pi} \int\left(\frac{\partial E_{x}}{\partial x_{0}}+\frac{\partial E_{y}}{\partial z_{0}}\right) \frac{e^{i k r}}{r} d x_{0} d y_{0},
\end{gathered}
$$

where $z_{1}>z_{0}>0 ; E_{x}\left(x_{0}, y_{0}, z_{0}\right)$ and $E_{y}\left(x_{0}, y_{0}, z_{0}\right)$ are given functions in the plane $z=z_{0}$. Formulae (1) present a general solution to the direct diffraction problem. They express the electric field $E$ at a point $x_{1}, y_{1}, z_{1}>z_{0}$ in terms of its boundary values $E_{x}\left(x_{0}, y_{0}, z_{0}\right), E_{y}\left(x_{0}, y_{0}, z_{0}\right)$ in a plane $z=z_{0}$.

Following the properly extended reasoning of Shewell and Wolf [2] we can obtain formulae for the inverse electromagnetic diffraction:
$E_{x}\left(x_{0}, y_{0}, z_{0}\right)=-\frac{1}{2 \pi} \int E_{x}\left(x_{1}, y_{1}, z_{1}\right) \frac{\partial}{\partial z_{1}}\left(\frac{e^{-i k r}}{r}\right) d x_{1} d y_{1}$,

[^0]$E_{y}\left(x_{0}, y_{0}, z_{0}\right)=-\frac{1}{2 \pi} \int E_{y}\left(x_{1}, y_{1}, z_{1}\right) \frac{\partial}{\partial z_{1}}\left(\frac{e^{-i k r}}{r}\right) d x_{1} d y_{1}$
\[

$$
\begin{equation*}
E_{z}\left(x_{0}, y_{0}, z_{0}\right)=-\frac{1}{2 \pi} \int\left(\frac{\partial E_{x}}{\partial x_{1}}+\frac{\partial E_{y}}{\partial y_{1}}\right) \frac{e^{-i k r}}{r} d x_{1} d y_{1} \tag{2}
\end{equation*}
$$

\]

Formulae (2) present evidently ( $0<z_{0}<z_{1}$ ) a general solution to the inverse diffraction problem. They express the electric field $E$ at a point $x_{0}, y_{0}, z_{0}$ in terms of its boundary values $E_{x}\left(x_{1}, y_{1}, z_{1}\right), E_{y}\left(x_{1}\right.$, $y_{1}, z_{1}$ ) in a plane $z=z_{1}$. Formulae (2) have been obtained under assumption [2] that the contribution introduced by the inhomogenous plane waves in the angular-spectrum representation [3, 4] of $E$ may be neglected.

## Aknowledgement

The author is greatly indebted to Prof. Albéric Boivin of the Laval University for many stimulating discussions. Thanks are also due the N.R.C and the Department of Physics of the Laval University for having made possible the author's stay at the Laval University.

## References

[1] Sherman G. C., J. Opt. Soc. Am. 57 (1965), 1490.
[2] Shewell R. J., Wolf E., J. Opt. Soc. Am. 58 (1968) 1596.
[3] Laror E., J. Opt. Soc. Am. 58 (1968), 1235.
[4] Laror E., J. Math. Phys. 9 (1968), 2001.
[5] Sherman G.C., Phys. Rev. Lett. 21 (1968), 761.
[6] Sherman G.C., J. Opt. Soc. Am. 59 (1969), 697.
[7] Jannson T., Karczewski B., Opt. Communications, 2, (1970), 201.
[8] Luneburg R. K., Mathematical Theory of Optics, Univ. of California Press, 1964.


[^0]:    *) Instytut Fizyki Politechniki Warszawskiej, Warszawa, ul. Koszykowa 75, Poland.

